

Neutrino decays over cosmological distances: GRBs

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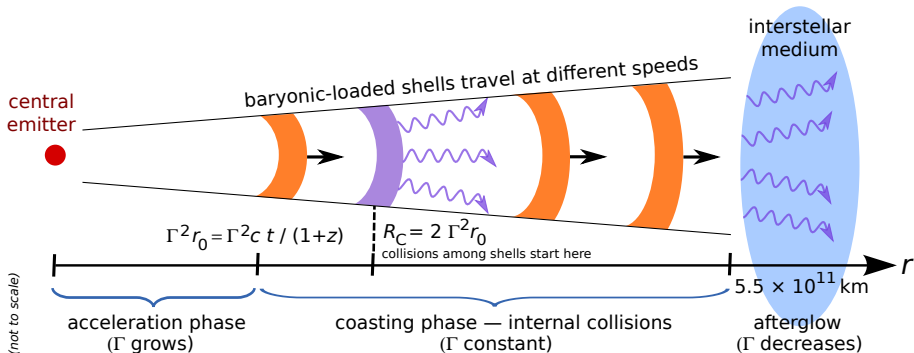


THE OHIO STATE UNIVERSITY

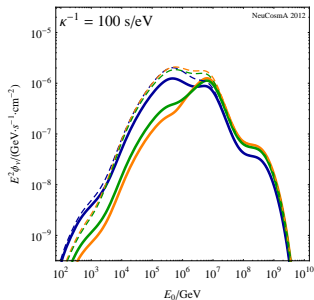
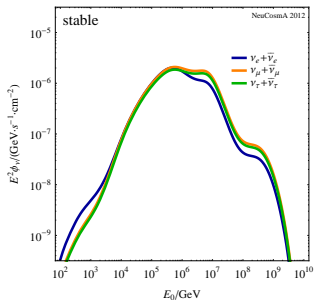


Fireball model: our current paradigm of how a GRB works

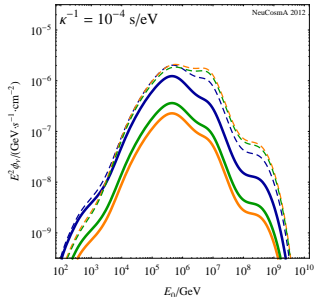
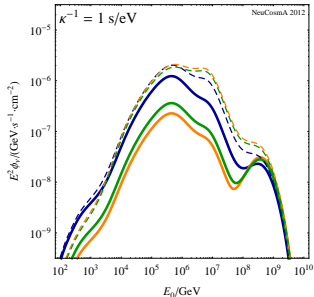
- relativistically-expanding blobs of plasma collide with each other and, in the process, emit UHE particles



UHE ν flux from a standard burst (ν_1 stable and $\nu_{2,3}$ unstable):

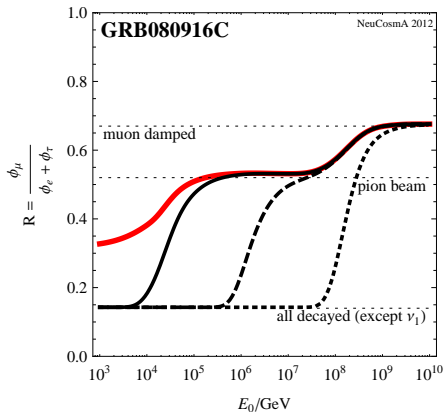
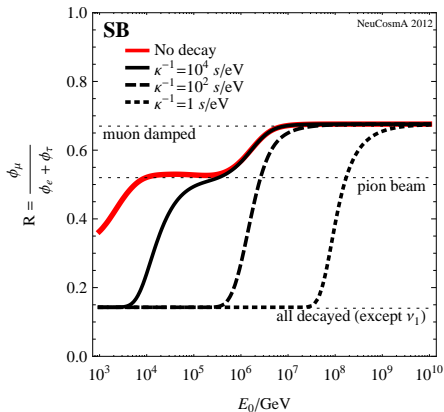


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Flavour ratio:

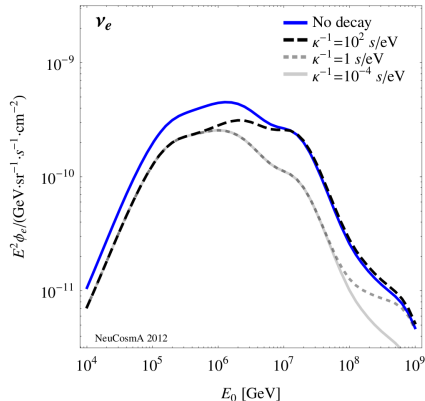
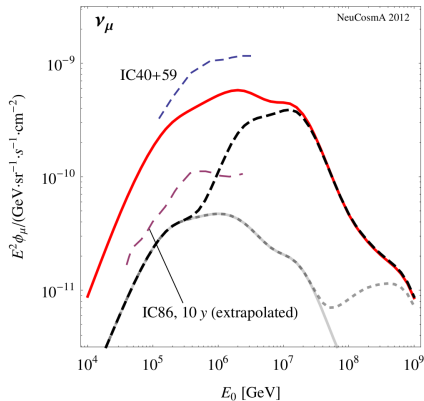
$$R = \frac{\phi_\mu}{\phi_e + \phi_\tau} \approx \frac{\text{number of muon tracks}}{\text{number of cascades}}$$



P. BAERWALD, MB, W. WINTER, *JCAP* **1210**, 020 (2012) [ARXIV:1208.4600]

Quasi-diffuse flux (stacking the same 117 GRBs from IC-40 analysis):

P. BAERWALD, MB, W. WINTER, *JCAP* **1210**, 020 (2012) [ARXIV:1208.4600]



Fewer muon tracks found because they decayed?

No reliable information on astrophysical neutrino sources can be obtained from muon tracks only

Backup slides

Joint production of UHECRs, ν 's, and γ 's:

power law $\sim E^{-\alpha p}$

broken power law

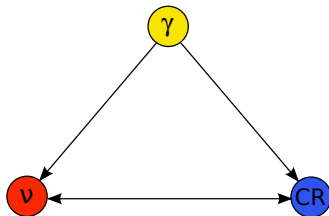
$$p \gamma \rightarrow \Delta^+ (1232) \rightarrow \begin{cases} n\pi^+, & \text{BR} = 1/3 \\ p\pi^0, & \text{BR} = 2/3 \end{cases}$$

$$\pi^+ \rightarrow \mu^+ \nu_\mu \rightarrow \bar{\nu}_\mu e^+ \nu_e \nu_\mu$$

$$\pi^0 \rightarrow \gamma\gamma$$

$$n \text{ (escapes)} \rightarrow p e^- \bar{\nu}_e$$

(Δ^+ : $\sim 50\%$ of all $p\gamma$ interactions)



After propagation, with flavour mixing:

$$\nu_e : \nu_\mu : \nu_\tau : p = 1 : 1 : 1 : 1$$

(“one ν_μ per cosmic ray”)

CR emission by n escape only is now strongly disfavoured

ICECUBE COLL., *Nature* **484**, 351 (2012)

AHLERS ET AL. *Astropart. Phys.* **35**, 87 (2011)

In a collision, UHE protons, photons, and neutrinos are emitted:

$$\underbrace{N'_p(E'_p)}_{\text{proton density at the source [GeV}^{-1} \text{ cm}^{-3}\text{]}} \quad \text{NeuCosmA} \quad \otimes \quad \underbrace{N'_\gamma(E'_\gamma)}_{\text{photon density at the source}}$$

$$= \underbrace{Q'_\nu(E'_\nu)}_{\text{ejected neutrino spectrum [GeV}^{-1} \text{ cm}^{-3} \text{ s}^{-1}\text{]}}$$

- ▶ From Fermi shock acceleration: $N'_p(E'_p) \propto E_p'^{-\alpha_p} e^{-E_p'/E_{p,\max}'}$
- ▶ Photon density at source has same shape as observed:

$$N'_\gamma(E'_\gamma) = \begin{cases} (E'_\gamma/E'_{\gamma,\text{break}})^{-\alpha_\gamma} & , E'_{\gamma,\text{min}} \leq E'_\gamma < E'_{\gamma,\text{break}} \\ (E'_\gamma/E'_{\gamma,\text{break}})^{-\beta_\gamma} & , E'_\gamma \geq E'_{\gamma,\text{break}} \\ 0 & , \text{otherwise} \end{cases}$$

$$\alpha_\gamma = 1, \beta_\gamma = 2.2, E'_{\gamma,\text{min}} = 0.2 \text{ eV}, E'_{\gamma,\text{break}} = 1 \text{ keV}$$

Normalise the densities at the source – for one collision:

► Photons:

$$\underbrace{\int E'_\gamma N'_\gamma (E'_\gamma) dE'_\gamma}_{\text{total energy density in photons}} = \frac{E_{\gamma\text{-sh}}^{\text{iso}}}{V'_{\text{iso}}}$$

baryonic loading (energy in p 's / energy in e 's + γ 's), e.g., 10

► Protons:

$$\underbrace{\int E'_p N'_p (E'_p) dE'_p}_{\text{total energy density in protons}} = \frac{1}{f_e} \frac{E_{\gamma\text{-sh}}^{\text{iso}}}{V'_{\text{iso}}}$$

NeuCosmA calculates the injected/ejected spectrum of secondaries (π , K , n , ν , etc.):

$$x \equiv E'/E'_p$$

$$y \equiv E'_p E'_\gamma / (m_p c^2)$$

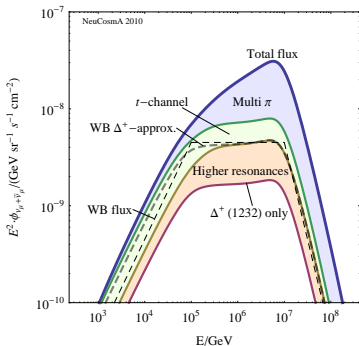
$$Q'(E') = \int_{E'}^{\infty} \frac{dE'_p}{E'_p} N'_p(E'_p) \int_0^{\infty} c dE'_\gamma N'_\gamma(E'_\gamma) R(x, y)$$

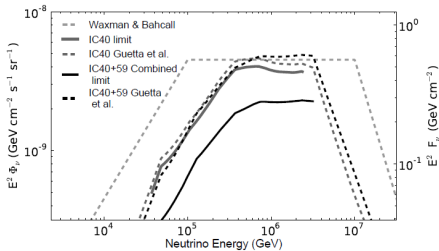
response function

R contains cross sections, multiplicities for different channels

What does NeuCosmA include?

- ▶ $p\gamma \rightarrow \Delta^+ (1232) \rightarrow \pi^0, \pi^+, \dots$
- ▶ extra K , n , π^- , multi- π production modes
- ▶ synchrotron losses of secondaries
- ▶ adiabatic cooling
- ▶ full photon spectrum
- ▶ neutrino flavour transitions





IceCube Collaboration:

- ▶ ν flux normalised to GRB γ fluence:

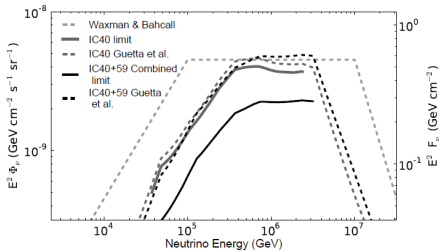
$$\int_0^\infty dE_\nu E_\nu F_\nu(E_\nu) \propto \int_{1 \text{ keV}}^{10 \text{ MeV}} d\varepsilon_\gamma \varepsilon_\gamma F_\gamma(\varepsilon_\gamma)$$

- ▶ quasi-diffuse ν flux from 117 GRBs
- ▶ **analytical calculation** – in tension with upper bounds

ICECUBE COLL., *Nature* **484**, 351 (2012)

AHLERS ET AL. *Astropart. Phys.* **35**, 87 (2011)

GUETTA ET AL. *Astropart. Phys.* **20**, 429 (2004)



More detailed particle physics (NeuCosmA):

- ▶ extra multi- π , K , n production modes
- ▶ synchrotron losses of secondaries
- ▶ adiabatic cooling
- ▶ full photon spectrum, etc.

ν flux \sim one order of magnitude lower

BAERWALD, HÜMMER, WINTER, *PRL* **108**, 231101 (2012)

See also: HE, LIU, WANG, *ApJ* **752**, 29 (2012)

IceCube Collaboration:

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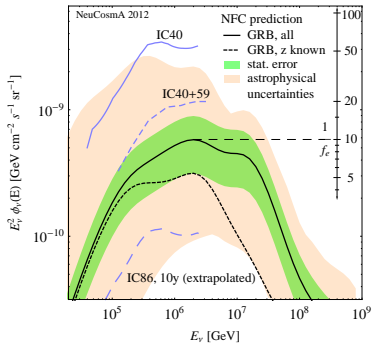
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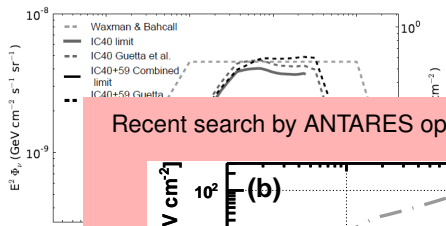
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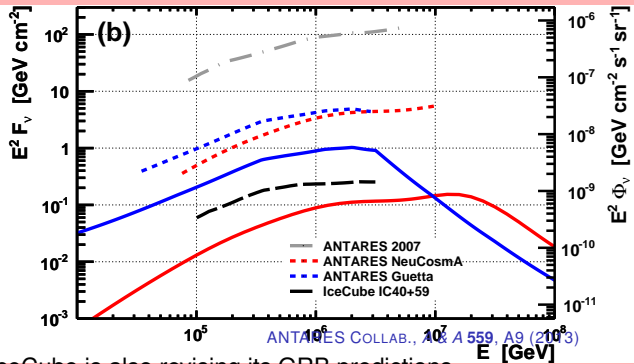




More detailed particle physics (NeuCosmA):

- ▶ extra multi- π , K , n production modes
- ▶ synchrotron losses of secondaries

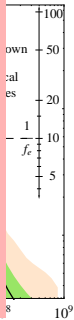
Recent search by ANTARES optimised for NeuCosmA:



ANTARES COLLAB., A & A 559, A9 (2013)

▶ IceCube is also revising its GRB predictions

wer
(2012)



E_ν [GeV]

- ▶ Same $n = 117$ GRBs, effective area, and parameters as used by the IC-40 analysis

- ▶ Calculate the associated neutrino flux for each burst and the stacked flux $F_\nu(E_\nu)$

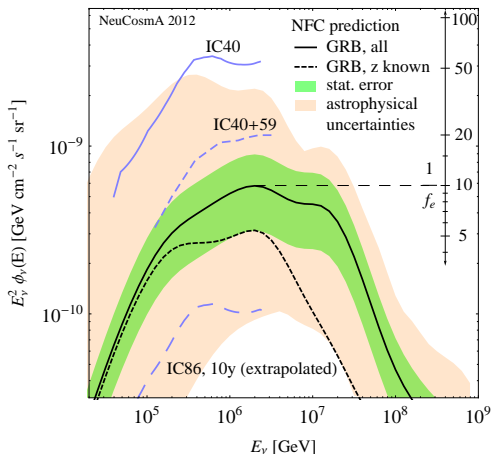
- ▶ Quasidiffuse flux:

$$\phi_\nu(E_\nu) = F_\nu(E_\nu) \frac{1}{4\pi} \frac{1}{n} \frac{667 \text{ bursts}}{\text{yr}}$$

- ▶ **Statistical uncertainty:** extrapolation of a few bursts to a quasidiffuse flux

- ▶ **Astrophysical uncertainty:**

- ▶ $0.001 \leq t_\nu [\text{s}] \leq 0.1$
- ▶ $200 \leq \Gamma \leq 500$
- ▶ $1.8 \leq \alpha_p \leq 2.2$
- ▶ $0.1 \leq \epsilon_e/\epsilon_B \leq 10$



S. HÜMMER, P. BAERWALD, AND W. WINTER,
Phys. Rev. Lett. **108**, 231101 (2012)